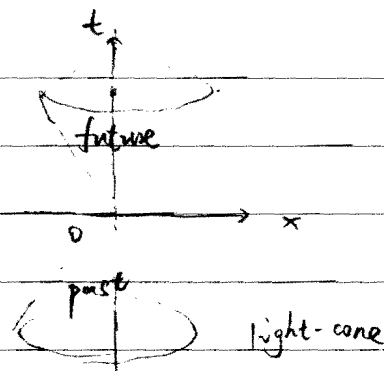


① retard potential.

$$x^\mu = (ct, x, y, z)$$

$$ds^2 = \eta_{\mu\nu} dx^\mu dx^\nu$$

$$\eta_{\mu\nu} = \begin{bmatrix} -1 & & & \\ & 1 & & \\ & & 1 & \\ & & & 1 \end{bmatrix}$$



null:  $ds^2 = 0$

$$\Rightarrow -c^2(t' - t)^2 + R^2 = 0$$

$$\Rightarrow t' = t - \frac{R}{c}$$

retard potential.

② Liénard-Wiechert potential

i) notation: Inertial reference frame I

Non-inertial reference frame  $\mathcal{N}^a$  (acceleration)

$\mathcal{N}^g$  (gravitation)

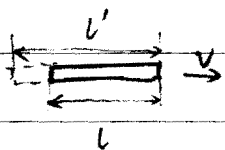
ii) Two facts.

a) The charge is Lorentz invariant

b) At every instant all physical quantities determined in non-inertial reference frame  $\mathcal{N}^a$  (or  $\mathcal{N}^g$ ) are equal to those determined in the co-moving (or local) inertial reference frame I.

iii) Origin of Liénard-Wiechert potential: an apparent change in the volume of a charge as seen at the observation point. (Griffiths, Intro. to Electrodynamics)

iv) Now consider a charge  $e = \rho V$  of length  $l$  and volume  $V$ . Assume



the charge constantly emits signal to "update" the field. (think of QED)

The potential at P at  $t$  in I is determined by all signals originating from different points of the charge at different moments and reaching P at the same moment  $t$ .

(But  $ds^2 = 0$  is always true for this condition, so we need retard time.)

For simplicity, consider the potential of a charge moving towards P.

$$\Delta t = \frac{l + v \Delta t}{c} \Rightarrow \Delta t = \frac{l/c}{1 - v/c}$$

The vector notation is easily obtained:

$$\Delta t = \frac{l/c}{1 - \frac{\vec{v} \cdot \vec{n}}{c}} \Rightarrow l' = c \Delta t = \frac{l}{1 - \frac{\vec{v} \cdot \vec{n}}{c}}$$

so  $V' = \frac{V}{1 - \frac{\vec{v} \cdot \vec{n}}{c}}$ , then the potential at P is

$$\varphi = \frac{1}{4\pi\epsilon_0} \cdot \frac{qV'}{r} = \frac{q}{4\pi\epsilon_0 r} \cdot \frac{1}{(1 - \frac{\vec{v} \cdot \vec{r}}{c^2})} \quad (\text{note: never forget retard time})$$

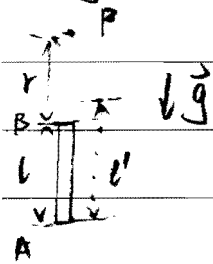
the vector Liénard-Wiechart potential vs

$$\vec{A} = \frac{q}{4\pi\epsilon_0 r c^2} \left[ \frac{\vec{v}}{1 - \frac{\vec{v} \cdot \vec{r}}{c^2}} \right]$$

The scalar and vector potential do not depend on the size of the size. so they are valid for a point charge.

② A rest charge in non-inertial reference frame

i) gravitational field  $N^g$



$$\varphi^g = \frac{1}{4\pi\epsilon_0} \cdot \frac{qV^g}{r^g}$$

Note: ~~P is in the local inertial reference frame~~ <sup>it doesn't matter</sup>

$$c_{AP}^g = c \left( 1 - \frac{g(r+l)}{2c^2} \right), \quad c_{BP}^g = c \left( 1 - \frac{gr}{2c^2} \right)$$

$$r^g = \frac{c_{BP}^g}{c} \cdot r, \quad \Delta t^g = \frac{r+l}{c_{AP}^g} - \frac{r}{c_{BP}^g}$$

$$l^g = c_{AP}^g \Delta t^g = r+l - \frac{c_{AP}^g}{c_{BP}^g} \cdot r \Rightarrow V^g = \frac{l^g}{l} \cdot V$$

$$\therefore \varphi^g = \frac{qV}{4\pi\epsilon_0 r} \cdot \frac{c}{c_{BP}^g} \cdot \frac{r+l - \frac{c_{AP}^g}{c_{BP}^g} \cdot r}{l} = \frac{qV}{4\pi\epsilon_0 r} \left( 1 + \frac{g \cdot r}{c^2} \right) + o(l)$$

point charge

the vector notation vs

$$\varphi^g = \frac{qV}{4\pi\epsilon_0 r} \left( 1 - \frac{\vec{g} \cdot \vec{r}}{c^2} \right)$$

ii) accelerated  $N^a$

the same procession  $\varphi^a = \frac{qV}{4\pi\epsilon_0 r} \left( 1 + \frac{\vec{a} \cdot \vec{r}}{c^2} \right)$

(note; no retarded time appears here)

$\vec{g} \leftrightarrow -\vec{a}$  Equivalence principle.

Ref. E. Fermi: Nuovo Cimento (1921).  $\varphi_{Fermi}^g = \frac{qV}{4\pi\epsilon_0 r} \left( 1 - \frac{\vec{g} \cdot \vec{r}}{2c^2} \right)$

But he get a factor  $\frac{1}{2}$  by some mistakes.

④ Charge moving in non-inertial reference frame.

We can just use  $V^g, r^g$  instead of  $V, r$  in the Liénard-Weichert potential.

$$\varphi^g = \frac{e}{4\pi\epsilon_0 r} \cdot \left[ \frac{1}{1 - \frac{\vec{v} \cdot \vec{n}}{c}} \cdot \left( 1 - \frac{\vec{g} \cdot \vec{r}}{c^2} \right) \right], \quad \vec{A}^g = \frac{e}{4\pi\epsilon_0 r} \cdot \left[ \frac{\vec{v}}{1 - \frac{\vec{v} \cdot \vec{n}}{c}} \cdot \left( 1 - \frac{\vec{g} \cdot \vec{r}}{c^2} \right) \right]$$

for  $\mathcal{N}^g$  and  $\vec{a} = -\vec{g}$  for  $\mathcal{N}^a$

Let the charge moving with  $\vec{a} = -\vec{g}$  in  $\mathcal{N}^g$  (but with  $\vec{v} = 0$ )

$$\text{then } \vec{E} = -\nabla\varphi^g - \frac{\partial \vec{A}^g}{\partial t} = \frac{e}{4\pi\epsilon_0} \left[ \left( \frac{\vec{n}}{r^2} + \frac{\vec{g} \cdot \vec{n}}{c^2 r} - \frac{\vec{g}}{c^2 r} \right) + \left( -\frac{\vec{g} \cdot \vec{n}}{c^2 r} \vec{n} + \frac{\vec{g}}{c^2 r} \right) \right]$$

$$= \frac{e}{4\pi\epsilon_0} \cdot \frac{\vec{n}}{r^2}, \text{ which reduce to Coulomb field.}$$

so No Radiation is obtained for a free falling charge.