

• 周期实信号：正交三角函数{coskω<sub>0</sub>t, sinkω<sub>0</sub>t}展开

exponential	triangular
$x(t) = \sum_{k=-\infty}^{\infty} C_k e^{jk\omega_0 t}$	$x(t) = \frac{a_0}{2} + \sum_{k=1}^{\infty} [a_k \cos(k\omega_0 t) + b_k \sin(k\omega_0 t)]$
$C_k = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jk\omega_0 t} dt$ (k = 0, ±1, ±2, ±3, ...)	$a_0 = \frac{2}{T} \int_0^T x(t) dt$ $a_k = \frac{2}{T} \int_0^T x(t) \cos(k\omega_0 t) dt$ $b_k = \frac{2}{T} \int_0^T x(t) \sin(k\omega_0 t) dt$ (k = 1, 2, 3, ...)

59

1

• 周期实信号

$$x(t) = \sum_{k=-\infty}^{\infty} C_k e^{jk\omega_0 t} \Rightarrow x^*(t) = \sum_{k=-\infty}^{\infty} C_k^* e^{-jk\omega_0 t}$$

以-k代替k:  $x^*(t) = \sum_{k=-\infty}^{\infty} C_{-k}^* e^{jk\omega_0 t}$

对实信号:  $x(t) = x^*(t) \Rightarrow x(t) = \sum_{k=-\infty}^{\infty} C_{-k}^* e^{jk\omega_0 t}$

比较  $x(t) = \sum_{k=-\infty}^{\infty} C_k e^{jk\omega_0 t} \Rightarrow C_k = C_{-k}^* \text{ or } C_{-k} = C_k^*$

$$\begin{aligned} \therefore x(t) &= \sum_{k=-\infty}^{\infty} C_k e^{jk\omega_0 t} = \sum_{k=-\infty}^{-1} C_k e^{jk\omega_0 t} + C_0 + \sum_{k=1}^{\infty} C_k e^{jk\omega_0 t} \\ &= \sum_{k=-\infty}^{-1} C_{-k}^* e^{jk\omega_0 t} + C_0 + \sum_{k=1}^{\infty} C_k e^{jk\omega_0 t} = \sum_{k=1}^{\infty} C_k^* e^{-jk\omega_0 t} + C_0 + \sum_{k=1}^{\infty} C_k e^{jk\omega_0 t} \\ &= C_0 + \sum_{k=1}^{\infty} (C_k e^{jk\omega_0 t} + C_k^* e^{-jk\omega_0 t}) \end{aligned}$$

59

2

令:  $C_k = |C_k| e^{j\theta_k}$

$$\begin{aligned} x(t) &= C_0 + \sum_{k=1}^{\infty} (C_k e^{jk\omega_0 t} + C_k^* e^{-jk\omega_0 t}) \\ &= C_0 + \sum_{k=1}^{\infty} (|C_k| e^{j\theta_k} e^{jk\omega_0 t} + |C_k| e^{-j\theta_k} e^{-jk\omega_0 t}) \\ &= C_0 + \sum_{k=1}^{\infty} |C_k| [e^{j(k\omega_0 t + \theta_k)} + e^{-j(k\omega_0 t + \theta_k)}] \\ &= C_0 + \sum_{k=1}^{\infty} 2|C_k| \cos(k\omega_0 t + \theta_k) \end{aligned}$$

C<sub>0</sub>是信号的直流分量

2|C<sub>k</sub>|是信号第k次谐波的振幅, θ<sub>k</sub>是信号第k次谐波的相位

• 任意周期实信号可由其直流分量和各次谐波分量合成

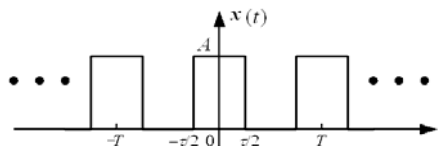
59

3

二、离散频谱

$$C_k = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jk\omega_0 t} dt \quad x(t) = \sum_{k=-\infty}^{\infty} C_k e^{jk\omega_0 t}, \quad |t| \leq T/2$$

例1 计算图示周期矩形脉冲信号的傅里叶级数



解:

$$C_k = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jk\omega_0 t} dt = \frac{1}{T} \int_{-\tau/2}^{\tau/2} A e^{-jk\omega_0 t} dt \quad \text{其中 } \omega_0 = \frac{2\pi}{T}$$

59

4

当k=0时,

$$C_0 = \frac{1}{T} \int_{-\tau/2}^{\tau/2} A e^0 dt = \frac{1}{T} \int_{-\tau/2}^{\tau/2} A dt = \frac{A\tau}{T}$$

当k≠0时,

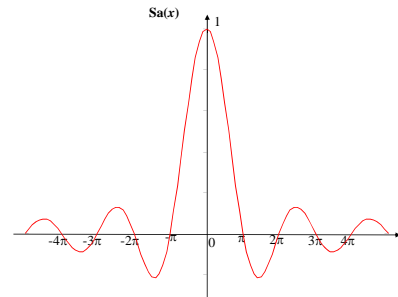
$$\begin{aligned} C_k &= \frac{1}{T} \int_{-\tau/2}^{\tau/2} A e^{-jk\omega_0 t} dt = \frac{A}{T - jk\omega_0} e^{-jk\omega_0 t} \Big|_{-\tau/2}^{\tau/2} = \frac{A}{T - jk\omega_0} \left( e^{-jk\omega_0 \tau/2} - e^{jk\omega_0 \tau/2} \right) \\ &= \frac{A}{T - jk\omega_0} \left[ -2j \sin\left(\frac{k\omega_0 \tau}{2}\right) \right] = \frac{A}{T} \frac{2 \sin(k\omega_0 \tau/2)}{k\omega_0} = \frac{A\tau}{T} \frac{\sin(k\omega_0 \tau/2)}{k\omega_0 \tau/2} \end{aligned}$$

定义: 抽样函数

$$\text{Sa}(x) = \frac{\sin x}{x}$$

59

5



$$C_0 = \frac{A\tau}{T}; \quad C_k = \frac{A\tau \sin(k\omega_0 \tau/2)}{T k\omega_0 \tau/2} = \frac{A\tau}{T} \text{Sa}(k\omega_0 \tau/2), \quad k \neq 0$$

$$\Rightarrow C_k = \frac{A\tau}{T} \text{Sa}(k\omega_0 \tau/2)$$

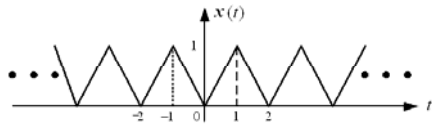
59

6

周期方波信号的傅里叶级数展开式为：

$$x(t) = \sum_{k=-\infty}^{\infty} C_k e^{jk\omega_0 t} = \frac{A\tau}{T} \sum_{k=-\infty}^{\infty} Sa(k\omega_0\tau/2) e^{jk\omega_0 t}, \quad |t| \leq \frac{T}{2}$$

**例2** 计算图示周期三角脉冲信号的傅里叶级数



**解：**

$$C_k = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} x(t) e^{-jk\omega_0 t} dt = \frac{1}{2} \left( \int_{-1}^0 -te^{-jk\omega_0 t} dt + \int_0^1 te^{-jk\omega_0 t} dt \right)$$

当k=0时，

$$C_0 = \frac{1}{2} \left( \int_{-1}^0 -te^0 dt + \int_0^1 te^0 dt \right) = \frac{1}{4} \left( -t^2 \Big|_{-1}^0 + t^2 \Big|_0^1 \right) = \frac{1}{2}$$

59

7

当k≠0时， $C_k = \frac{1}{2jk\omega_0} (te^{-jk\omega_0 t} \Big|_{-1}^0 - \int_{-1}^0 e^{-jk\omega_0 t} dt - te^{-jk\omega_0 t} \Big|_0^1 + \int_0^1 e^{-jk\omega_0 t} dt)$

$$\therefore \omega_0 = \frac{2\pi}{T} = \frac{2\pi}{2} = \pi$$

$$\therefore C_k = \frac{1}{2jk\pi} (0 + e^{jk\pi} - \int_{-1}^0 e^{-jk\pi} dt - e^{-jk\pi} + 0 + \int_0^1 e^{-jk\pi} dt)$$

$$= \frac{1}{2jk\pi} \left( 2j \sin k\pi - \frac{1}{-jk\pi} e^{-jk\pi} \Big|_{-1}^0 + \frac{1}{-jk\pi} e^{-jk\pi} \Big|_0^1 \right)$$

$$= \frac{1}{2jk\pi} \left( -\frac{1}{-jk\pi} + \frac{1}{-jk\pi} e^{jk\pi} + \frac{1}{-jk\pi} e^{-jk\pi} - \frac{1}{-jk\pi} \right)$$

$$= \frac{1}{2(k\pi)^2} (-1 + e^{jk\pi} + e^{-jk\pi} - 1) = \frac{1}{2(k\pi)^2} (2 \cos k\pi - 2)$$

$$= \frac{1}{(k\pi)^2} (\cos k\pi - 1) = \begin{cases} 0 & k \text{ 为偶数} \\ -\frac{2}{(k\pi)^2} & k \text{ 为奇数} \end{cases}$$

59

8

$$\therefore C_k = \begin{cases} -2/(k\pi)^2 & k \text{ 为奇数} \\ 1/2 & k = 0 \\ 0 & k \text{ 为非0偶数} \end{cases}$$

周期三角脉冲信号的傅里叶级数展开式为：

$$x(t) = \sum_{k=-\infty}^{\infty} C_k e^{jk\omega_0 t} = \frac{1}{2} - \sum_{m=-\infty}^{\infty} \frac{2}{[(2m-1)\pi]^2} e^{j(2m-1)\pi t}, \quad |t| \leq T/2 = 1$$

• 周期信号x(t)可以分解为不同频率虚指数信号之和

$$x_T(t) = \sum_{k=-\infty}^{\infty} C_k e^{jk\omega_0 t}$$

$C_k$ 是频率的函数，反映了组成信号各谐波的幅度和相位随频率变化的规律，称为**频谱函数**

59

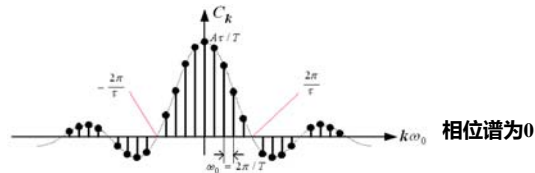
9

$$C_k = |C_k| e^{j\theta_k}$$

振幅谱

相位谱

- 周期信号的**频谱图**：信号各次谐波对应的 $C_k$ 线状分布图
- 周期矩形脉冲信号的**频谱图**  $C_k = \frac{A\tau}{T} Sa(k\omega_0\tau/2)$



59

10

• 相位谱的作用



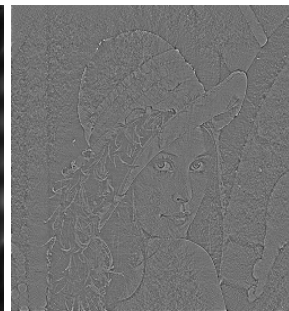
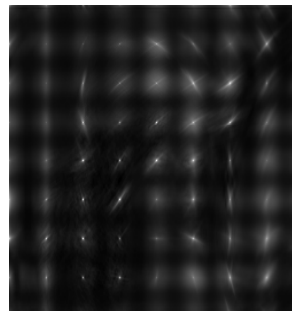
59

11

二维傅里叶级数

幅频不变，零相位

幅频为常数，相位不变



59

12

- 周期信号频谱的三个特点：**离散性、谐波性、收敛性**
  - ◆ 频谱由不连续的线条组成，每一条线代表一个谐波分量，称为**不连续频谱或离散频谱**
  - ◆ 每条线谱只出现在基波频率 $\omega_0$ 的整数倍的频率上，频谱中不可能存在任何具有频率为基波频率非整数倍的分量
  - ◆ 各条线谱的高度(即各次谐波的幅值)，随着谐波次数 $k$ 的增大而逐渐减少的，直至零

**信号周期 $T$ 越大， $\omega_0$ 就越小，则谱线越密**  
**信号周期 $T$ 越小， $\omega_0$ 就越大，则谱线越疏**

**信号时域波形变化越平缓，高次谐波成分就越少，幅度频谱衰减越快**  
**信号时域波形跳变越多，高次谐波成分就越多，幅度频谱衰减越慢**

59 13

### 三、离散频谱的基本性质

#### 1. 线性特性

若  $x_1(t) \rightarrow C_{1k}, x_2(t) \rightarrow C_{2k}$

则有  $a_1 \cdot x_1(t) + a_2 \cdot x_2(t) \rightarrow a_1 \cdot C_{1k} + a_2 \cdot C_{2k}$

**证:**

$$C_k = \frac{1}{T} \int_{-T/2}^{T/2} [a_1 \cdot x_1(t) + a_2 \cdot x_2(t)] e^{-jk\omega_0 t} dt$$

$$= a_1 \frac{1}{T} \int_{-T/2}^{T/2} x_1(t) e^{-jk\omega_0 t} dt + a_2 \frac{1}{T} \int_{-T/2}^{T/2} x_2(t) e^{-jk\omega_0 t} dt$$

$$= a_1 \cdot C_{1k} + a_2 \cdot C_{2k}$$

#### 2. 对称特性

**(1) 若 $x(t)$ 为实信号，则  $|C_k| = |C_{-k}|, \theta_k = -\theta_{-k}$**

59 14

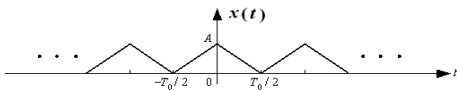
**证:**  $C_k = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jk\omega_0 t} dt, C_{-k} = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{jk\omega_0 t} dt$

$x(t)$ 为实信号  $\Rightarrow x(t) = x^*(t)$

$$C_{-k} = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{jk\omega_0 t} dt = \left[ \frac{1}{T} \int_{-T/2}^{T/2} x^*(t) e^{-jk\omega_0 t} dt \right]^* = \left[ \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jk\omega_0 t} dt \right]^* = C_k^*$$

$\therefore |C_{-k}| = |C_k^*| = |C_k|, \theta_{-k} = \arctan \frac{\text{Im} C_{-k}}{\text{Re} C_{-k}} = \arctan \frac{\text{Im} C_k^*}{\text{Re} C_k^*} = \arctan \frac{-\text{Im} C_k}{\text{Re} C_k} = -\theta_k$

**(2) 若 $x(t)$ 为实偶信号:  $x(t) = x(-t)$**



**$C_k$ 是 $k\omega_0$ 的实偶函数  $\Rightarrow$  振幅谱偶对称; 相位谱为0**

59 15

**证:**  $x(t)$ 为实信号, 则  $C_k^* = C_{-k}$

$x(t)$ 为偶信号  $\Rightarrow x(t) = x(-t)$

$$C_{-k} = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-j(-k)\omega_0 t} dt = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jk\omega_0(-t)} dt = \frac{1}{T} \int_{-T/2}^{T/2} x(-t') e^{-jk\omega_0 t'} (-dt')$$

$$= \frac{1}{T} \int_{-T/2}^{T/2} x(-t') e^{-jk\omega_0 t'} dt' = \frac{1}{T} \int_{-T/2}^{T/2} x(t') e^{-jk\omega_0 t'} dt' = C_k$$

**令:  $C_k = \alpha_k + j\beta_k$**

$$\alpha_k + j\beta_{-k} = C_{-k} = C_k = \alpha_k + j\beta_k \Rightarrow \alpha_k = \alpha_{-k}, \beta_k = \beta_{-k}$$

$$\alpha_{-k} + j\beta_{-k} = C_{-k} = C_k^* = \alpha_k - j\beta_k \Rightarrow \alpha_k = \alpha_{-k}, \beta_k = -\beta_{-k}$$

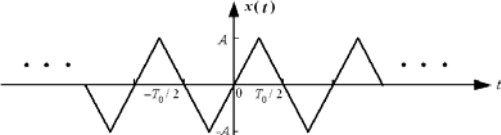
$$\beta_k = \beta_{-k} = -\beta_{-k} \Rightarrow \beta_{-k} = 0 \Rightarrow \beta_k = 0$$

$\therefore C_k = \alpha_k = \alpha_{-k} = C_{-k}$

**$C_k$ 是 $k\omega_0$ 的实偶函数**

59 16

**(3) 若 $x(t)$ 为实奇信号:  $x(t) = -x(-t)$**



**$C_k$ 是 $k\omega_0$ 的纯虚奇函数  $\Rightarrow$  振幅谱偶对称; 相位谱奇对称**

**证:**  $x(t)$ 为实信号, 则  $C_k^* = C_{-k}$

$x(t)$ 为奇信号  $\Rightarrow x(t) = -x(-t)$  类似  $\Rightarrow C_{-k} = -C_k$

$$\alpha_{-k} + j\beta_{-k} = C_{-k} = -C_k = -\alpha_k - j\beta_k \Rightarrow \alpha_k = -\alpha_{-k}, \beta_k = -\beta_{-k}$$

$$\alpha_{-k} + j\beta_{-k} = C_{-k} = C_k^* = \alpha_k - j\beta_k \Rightarrow \alpha_k = \alpha_{-k}, \beta_k = -\beta_{-k}$$

59 17

$$\alpha_k = \alpha_{-k} = -\alpha_{-k} \Rightarrow \alpha_{-k} = 0 \Rightarrow \alpha_k = 0$$

$$\therefore C_k = j\beta_k = -j\beta_{-k} = -C_{-k}$$

**$C_k$ 是 $k\omega_0$ 的纯虚奇函数**

$\Rightarrow |C_k| = |j\beta_k| = |-j\beta_{-k}| = |j\beta_{-k}| = |C_{-k}|$

if  $\beta_k > 0$ , then  $\beta_{-k} < 0$

$$\Rightarrow C_k = j\beta_k = \beta_k e^{j\frac{\pi}{2}}, C_{-k} = j\beta_{-k} = -j|\beta_{-k}| = |\beta_{-k}| e^{-j\frac{\pi}{2}}$$

if  $\beta_k < 0$ , then  $\beta_{-k} > 0$

$$\Rightarrow C_{-k} = j\beta_{-k} = \beta_{-k} e^{j\frac{\pi}{2}}, C_k = j\beta_k = -j|\beta_k| = |\beta_k| e^{-j\frac{\pi}{2}}$$

59 18

### 3. 周期信号功率与离散频谱的关系

#### 帕塞瓦尔(Parseval)功率守恒定理

$$P = \frac{1}{T} \int_{-T/2}^{T/2} |x(t)|^2 dt = \sum_{k=-\infty}^{\infty} |C_k|^2$$

物理意义:

周期信号平均功率按各谐波成分的振幅大小分配给各分量

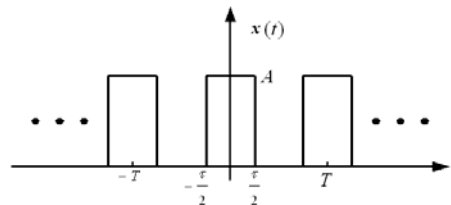
证:  $x(t) = \sum_{k=-\infty}^{\infty} C_k e^{jk\omega_0 t}, |t| \leq T/2$

$$\begin{aligned} P &= \frac{1}{T} \int_{-T/2}^{T/2} |x(t)|^2 dt = \frac{1}{T} \int_{-T/2}^{T/2} x(t)x^*(t) dt = \frac{1}{T} \int_{-T/2}^{T/2} \left[ \sum_{k=-\infty}^{\infty} C_k e^{jk\omega_0 t} \right] \left[ \sum_{n=-\infty}^{\infty} C_n^* e^{-jn\omega_0 t} \right] dt \\ &= \frac{1}{T} \int_{-T/2}^{T/2} \left[ \sum_{k=-\infty}^{\infty} C_k e^{jk\omega_0 t} \right] \left[ \sum_{n=-\infty}^{\infty} C_n^* e^{-jn\omega_0 t} \right] dt = \sum_{k=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} C_k C_n^* \frac{1}{T} \int_{-T/2}^{T/2} e^{j(k-n)\omega_0 t} dt \\ &= \sum_{k=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} C_k C_n^* \frac{1}{T} \int_{-T/2}^{T/2} e^{j(k-n)\omega_0 t} dt = \sum_{k=-\infty}^{\infty} C_k C_k^* = \sum_{k=-\infty}^{\infty} |C_k|^2 \end{aligned}$$

59

19

例3 求周期矩形脉冲信号在其有效带宽(0~2π/τ)内谐波分量所具有的平均功率占整个信号平均功率的百分比(其中A=1, T=1/4, τ=1/20)



解: 周期矩形脉冲的傅里叶级数为:

$$C_k = \frac{A\tau}{T} \text{Sa}(k\omega_0\tau/2) = \frac{1 \times 1/20}{1/4} \text{Sa}\left(k \frac{2\pi}{1/4} \cdot \frac{1}{20} \cdot \frac{1}{2}\right) = \frac{1}{5} \text{Sa}\left(\frac{k\pi}{5}\right)$$

59

20

信号的平均功率为:

$$P = \frac{1}{T} \int_{-T/2}^{T/2} |x(t)|^2 dt = \frac{1}{T} A^2 \tau = \frac{1/20}{1/4} = \frac{1}{5}$$

$$k\omega_0 = k \frac{2\pi}{T} \leq \frac{2\pi}{\tau} \Rightarrow k \leq \frac{T}{\tau} = \frac{1/4}{1/20} = 5$$

包含在有效带宽(0~2π/τ)内的各谐波平均功率为:

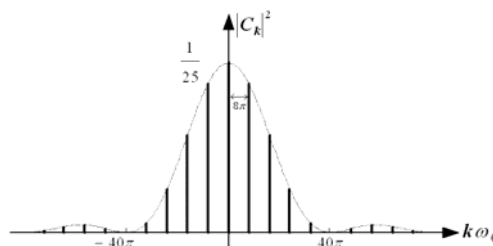
$$P_1 = \sum_{k=-5}^5 |C_k|^2 = \sum_{k=-5}^5 \left[ \frac{1}{5} \text{Sa}\left(\frac{k\pi}{5}\right) \right]^2 \approx 0.1806$$

$$\therefore \frac{P_1}{P} = \frac{0.1806}{0.2000} = 90\%$$

59

21

$$|C_k|^2 = \frac{1}{25} [\text{Sa}(k\pi/5)]^2, \quad \omega_0 = \frac{2\pi}{T} = \frac{2\pi}{1/4} = 8\pi$$



59

22

### 4. 截断离散频谱的高频成分对x(t)引入的失真

(1)有限项截断后, 傅里叶级数于最小均方误差意义逼近原信号均方误差

$$\begin{aligned} E &= \int_{-T}^T |e(t)|^2 dt = \int_{-T}^T [x(t) - x_1(t)][x^*(t) - x_1^*(t)] dt = \int_{-T}^T x(t)x^*(t) dt - \int_{-T}^T x(t) \sum_{k=-N}^N C_k e^{jk\omega_0 t} dt \\ &\quad - \int_{-T}^T x_1^*(t) \sum_{k=-N}^N C_k^* e^{-jk\omega_0 t} dt + \int_{-T}^T \sum_{k=-N}^N \sum_{n=-N}^N C_k^* C_n e^{j(n-k)\omega_0 t} dt = \int_{-T}^T |x(t)|^2 dt - \sum_{k=-N}^N C_k \int_{-T}^T x^*(t) e^{jk\omega_0 t} dt \\ &\quad - \sum_{k=-N}^N C_k^* \int_{-T}^T x(t) e^{-jk\omega_0 t} dt + \sum_{k=-N}^N \sum_{n=-N}^N C_k^* C_n \int_{-T}^T e^{j(n-k)\omega_0 t} dt = \int_{-T}^T |x(t)|^2 dt - \sum_{k=-N}^N C_k \left[ \int_{-T}^T x(t) e^{-jk\omega_0 t} dt \right]^* \\ &\quad - \sum_{k=-N}^N C_k^* C_k T + \sum_{k=-N}^N \sum_{n=-N}^N C_k^* C_n T = \int_{-T}^T |x(t)|^2 dt - \sum_{k=-N}^N C_k C_k^* T = \int_{-T}^T |x(t)|^2 dt - T \sum_{k=-N}^N |C_k|^2 \\ &= T \sum_{k=-\infty}^{\infty} |C_k|^2 - T \sum_{k=-N}^N |C_k|^2 = T \sum_{k=-\infty}^{-(N+1)} |C_k|^2 + T \sum_{k=N+1}^{\infty} |C_k|^2 \end{aligned}$$

当N→∞, E单调趋于0

59

23

### (2)吉布斯现象

用有限次谐波分量来近似原信号, 在x(t)的不连续点将出现过冲, 过冲峰值不随N增加而减少, 约为跳变值的9%

以方波为例: A=1, T=2, τ=T/2=1

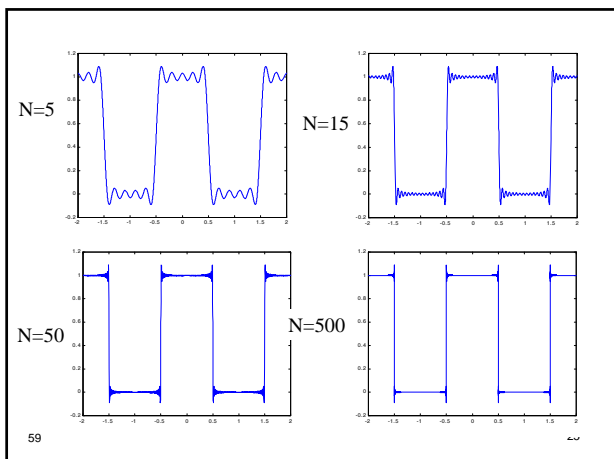
$$C_k = \frac{A\tau}{T} \text{Sa}(k\omega_0\tau/2) = \frac{1 \times 1}{2} \text{Sa}\left(k \frac{2\pi}{2} \times \frac{1}{2}\right) = \frac{1}{2} \text{Sa}(k\pi/2)$$

$$C_0 = 1/2; \quad C_{2n} = 0, n \neq 0; \quad C_{2n+1} = \frac{\sin[(2n+1)\pi/2]}{(2n+1)\pi}, \quad C_{-2n+1} = C_{-(2n+1)}$$

$$\begin{aligned} \therefore x(t) &= \sum_{k=-\infty}^{\infty} C_k e^{jk\omega_0 t} = C_0 + \sum_{n=0}^{\infty} [C_{2n+1} e^{j(2n+1)\pi t} + C_{-(2n+1)} e^{-j(2n+1)\pi t}] \\ &= \frac{1}{2} + \sum_{n=0}^{\infty} C_{2n+1} [e^{j(2n+1)\pi t} + e^{-j(2n+1)\pi t}] = \frac{1}{2} + 2 \sum_{n=0}^{\infty} \frac{\sin[(2n+1)\pi/2]}{(2n+1)\pi} \cos[(2n+1)\pi t] \\ &= \frac{1}{2} + 2 \sum_{n=0}^{\infty} \frac{(-1)^n}{(2n+1)\pi} \cos[(2n+1)\pi t] = \frac{1}{2} + \frac{2}{\pi} \cos(\pi t) - \frac{2}{3\pi} \cos(3\pi t) + \dots \end{aligned}$$

59

24



§ 1.2 能量信号的频谱分析

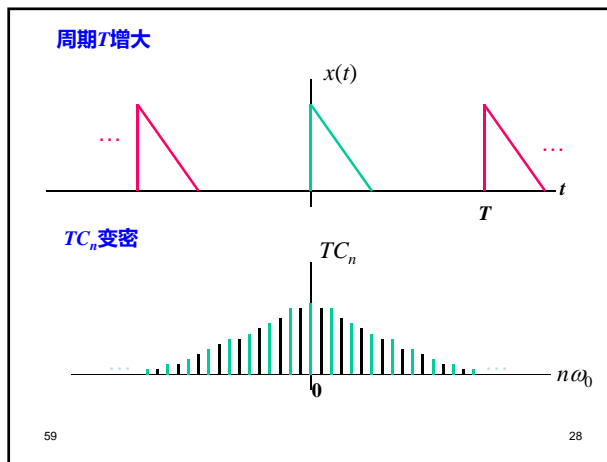
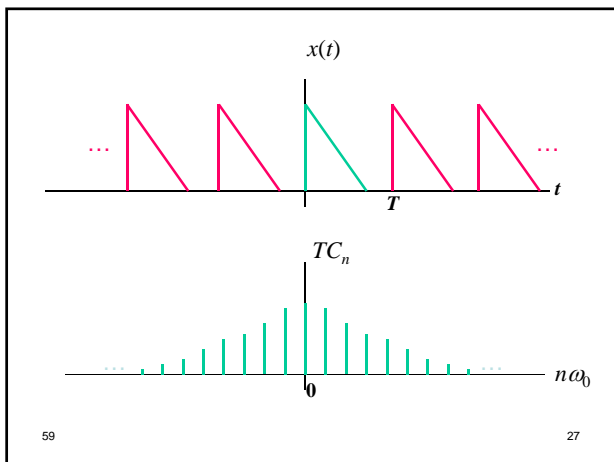
一、傅里叶变换(Fourier transform)

1. 从傅里叶级数到傅里叶变换

- 讨论周期T增加对离散频谱的影响

周期为T 宽度为τ的周期矩形脉冲的Fourier级数为：

$$C_n = \frac{A\tau}{T} \text{Sa}(n\omega_0\tau/2) \Rightarrow TC_n = A\tau \text{Sa}(n\omega_0\tau/2)$$



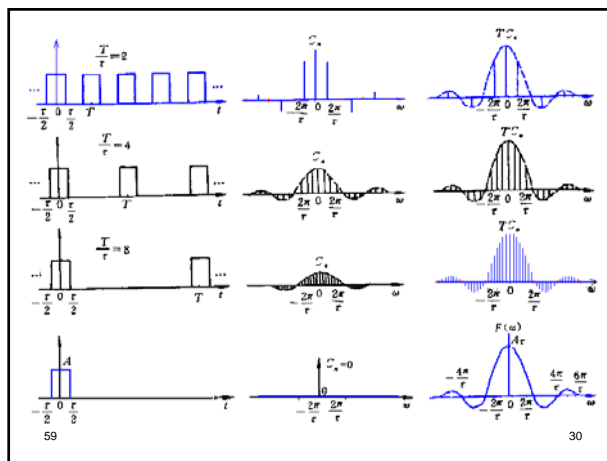
• 周期矩形脉冲信号的频谱

$$x(t) = \begin{cases} A & |t| \leq \frac{\tau}{2} \\ 0 & \frac{\tau}{2} < |t| \leq \frac{T}{2} \end{cases}$$

傅里叶级数展开式为：

$$x(t) = \sum_{n=-\infty}^{\infty} C_n e^{jn\omega_0 t}, |t| \leq T/2$$


傅里叶级数为：

$$C_n = \frac{A\tau}{T} \frac{\sin(n\omega_0\tau/2)}{(n\omega_0\tau/2)} = \frac{A\tau}{T} \text{Sa}(n\omega_0\tau/2), \quad n = 0, \pm 1, \pm 2, \dots$$


当  $T \rightarrow +\infty, \omega_0 = \frac{2\pi}{T} \rightarrow 0$

信号  $x(t)$  从周期信号变成非周期信号

令  $\omega_n = n\omega_0, \rightarrow \Delta\omega_n = \omega_n - \omega_{n-1} = \omega_0 \therefore T = \frac{2\pi}{\Delta\omega_n}$



周期信号:  $C_n = \frac{1}{T} \int x(t) e^{-jn\omega_0 t} dt \quad x(t) = \sum_{n=-\infty}^{\infty} C_n e^{jn\omega_0 t}, |t| \leq T/2$

非周期信号:  $x(t) = \lim_{\Delta\omega_n \rightarrow 0} \frac{1}{2\pi} \sum_{n=-\infty}^{+\infty} (TC_n) e^{j\omega_n t} \Delta\omega_n$

$\frac{1}{T} = \frac{\Delta\omega_n}{2\pi}$

$x(t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} X(\omega) e^{j\omega t} d\omega \quad \lim_{T \rightarrow \infty} TC_n = X(\omega)$

周期信号:  $C_n = \frac{1}{T} \int x(t) e^{-jn\omega_0 t} dt \quad x(t) = \sum_{n=-\infty}^{\infty} C_n e^{jn\omega_0 t}, |t| \leq T/2$

非周期信号:

$X(\omega) = \lim_{\Delta\omega_n \rightarrow 0} TC_n = \lim_{\Delta\omega_n \rightarrow 0} T \cdot \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jn\omega_0 t} dt = \lim_{\Delta\omega_n \rightarrow 0} \int_{-T/2}^{T/2} x(t) e^{-j\omega_n t} dt$

$X(\omega) = \int_{-\infty}^{+\infty} x(t) e^{-j\omega t} dt$

---

2. 傅里叶变换对

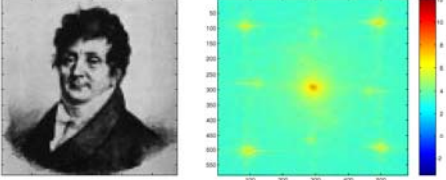
傅里叶正变换:

$X(\omega) = \int_{-\infty}^{+\infty} x(t) e^{-j\omega t} dt$  或  $X(f) = \int_{-\infty}^{\infty} x(t) e^{-j2\pi f t} dt$

傅里叶反变换:

$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\omega) e^{j\omega t} d\omega$  或  $x(t) = \int_{-\infty}^{\infty} X(f) e^{j2\pi f t} df$

两维Fourier transform



Fourier      Fourier transformed

582 by 582 pixels      图像两维傅里叶变换幅度的对数

- 符号表示:
 
$$X(\omega) = F[x(t)], \quad x(t) = F^{-1}[X(\omega)]$$
 或  $x(t) \xleftrightarrow{F} X(\omega)$
- 傅里叶反变换:
 
$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\omega) e^{j\omega t} d\omega$$

物理意义: 非周期信号可以分解为无数个频率为  $\omega$ 、复振幅为  $[X(\omega)/2\pi]d\omega$  的复指数信号  $e^{j\omega t}$  的线性组合

---

- 称信号  $x(t)$  的傅里叶变换  $X(\omega)$  为其**连续频谱**
- 一个信号满足狄里赫利条件, 其傅里叶变换存在

狄里赫利条件

- 非周期信号在无限区间上绝对可积
 
$$\int_{-\infty}^{\infty} |x(t)| dt < \infty$$
 充分不必要条件
- 在任意有限区间内, 信号只有有限个最大值和最小值
 必要不充分条件
- 在任意有限区间内, 信号仅有有限个不连续点, 且这些点必须是有限值
 必要不充分条件

---

一个连续信号  $x(t)$  的傅里叶变换  $X(\omega)$  是唯一的

二、连续频谱

- 周期信号的频谱为**离散频谱**
- 非周期信号的频谱为**连续频谱**
- 周期信号的频谱为  $C_n$  的分布, 表示每个谐波分量的复振幅
- 非周期信号的频谱为  $TC_n$  的分布, 表示每单位带宽内所有谐波分量合成的复振幅(即**频谱密度函数**)

两者关系:

$$X(\omega) = \lim_{T \rightarrow \infty} TC_n$$

$$C_n = \frac{X(\omega)}{T} \Big|_{\omega=n\omega_0}$$


---

- 连续频谱密度函数(简称**连续频谱**):

$$X(\omega) = \int_{-\infty}^{+\infty} x(t) e^{-j\omega t} dt$$

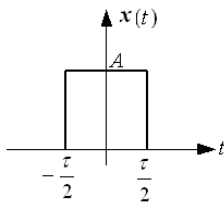
$X(\omega) = |X(\omega)| e^{j\varphi(\omega)} = R(\omega) + jI(\omega)$

振幅密度谱  
简称振幅谱

相位密度谱  
简称相位谱

---

**例1 计算图示非周期矩形脉冲信号的频谱函数**



**解:**  $x(t) = \begin{cases} A, & |t| \leq \tau/2 \\ 0, & |t| > \tau/2 \end{cases}$

可记作:  $x(t) = A \text{rect}\left(\frac{t}{\tau}\right)$

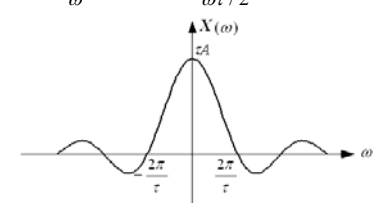
类似, 幅度A、中心 $t_0$ 、宽度 $\tau$ 的矩形脉冲可记作:

$$x(t) = A \text{rect}\left(\frac{t-t_0}{\tau}\right)$$

由傅里叶正变换定义式, 可得:

$$X(\omega) = \int_{-\infty}^{\infty} x(t) e^{-j\omega t} dt = \int_{-\tau/2}^{\tau/2} A \cdot e^{-j\omega t} dt = \frac{A}{-j\omega} e^{-j\omega t} \Big|_{-\tau/2}^{\tau/2}$$

$$= \frac{A}{-j\omega} (e^{-j\omega\tau/2} - e^{j\omega\tau/2}) = \frac{A}{-j\omega} [-2j \sin(\omega\tau/2)]$$

$$= \frac{2A \sin(\omega\tau/2)}{\omega} = A\tau \frac{\sin(\omega\tau/2)}{\omega\tau/2} = A\tau \cdot \text{Sa}(\omega\tau/2)$$


分析:

- 非周期矩形脉冲信号的频谱是**连续频谱**, 其形状与周期矩形脉冲信号离散频谱的包络线相似
- 周期信号的**离散频谱**可通过对非周期信号的连续频谱等间隔取样求得
- 信号在**时域有限**, 则在**频域将无限延续**
- 信号的频谱分量主要集中在**零频到第一个过零点**之间, 工程中往往将此宽度作为**有效带宽**
- 脉冲宽度 $\tau$ 越窄, 有限带宽越宽, 高频分量越多即信号信息量大、传输速度快, 传送信号所占用的频带越宽

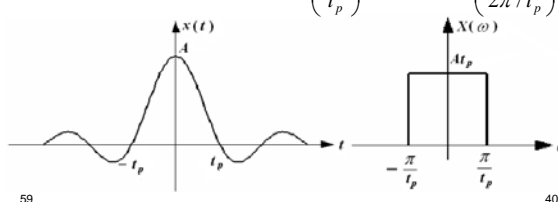
- 时域抽样函数**: 主瓣幅度A、中心 $t_0$ 、第一零点和中心距离 $t_p$

$$x(t) = A \text{Sa}\left[\frac{\pi(t-t_0)}{t_p}\right]$$

若 $t_0=0$ :

$$x(t) = A \text{Sa}\left(\frac{\pi t}{t_p}\right)$$

- 利用傅里叶反变换可证明:  $A \text{Sa}\left(\frac{\pi t}{t_p}\right) \xleftrightarrow{F} A t_p \text{rect}\left(\frac{\omega}{2\pi/t_p}\right)$



**证:**  $x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\omega) e^{j\omega t} d\omega = \frac{1}{2\pi} \int_{-\pi/t_p}^{\pi/t_p} A t_p \cdot e^{j\omega t} d\omega = \frac{1}{2\pi} \frac{A t_p}{jt} e^{j\omega t} \Big|_{-\pi/t_p}^{\pi/t_p}$

$$= \frac{A t_p}{2\pi jt} (e^{j\pi t/t_p} - e^{-j\pi t/t_p}) = \frac{A t_p}{2\pi jt} [2j \sin(\pi t/t_p)]$$

$$= \frac{A t_p \sin(\pi t/t_p)}{\pi t} = A \frac{\sin(\pi t/t_p)}{\pi t/t_p} = A \text{Sa}(\pi t/t_p)$$


---

**例2 计算指数信号 $e^{-at}$ ( $a>0$ )的频谱函数**

**解:** 由傅里叶正变换定义式, 可得:

$$X(\omega) = \int_{-\infty}^{\infty} x(t) e^{-j\omega t} dt = \int_{-\infty}^0 e^{-at} \cdot e^{-j\omega t} dt + \int_0^{\infty} e^{-at} \cdot e^{-j\omega t} dt$$

$$= \frac{1}{a-j\omega} e^{-(a-j\omega)t} \Big|_{-\infty}^0 + \frac{1}{-(a+j\omega)} e^{-(a+j\omega)t} \Big|_0^{\infty} = \frac{1}{a-j\omega} + \frac{1}{a+j\omega} = \frac{2a}{a^2 + \omega^2}$$

$\therefore |X(\omega)| = \frac{2a}{a^2 + \omega^2}, \quad \varphi(\omega) = 0$

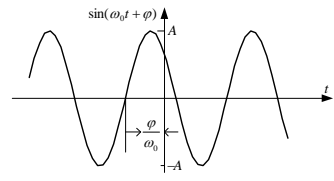
### § 1.3 信号的基本概念

#### 一、典型信号

##### 1. 正弦信号

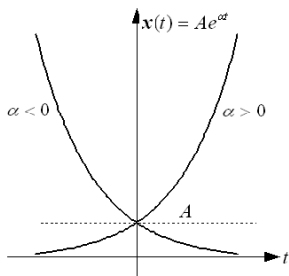
$$x(t) = A \sin(\omega_0 t + \varphi)$$

A: 振幅  $\omega_0$ : 角频率(弧度/秒)  $\varphi$ : 初始相位



### 2. 实指数信号

$$x(t) = Ae^{\alpha t}$$



59

43

### 3. 虚指数信号

$$x(t) = e^{j\omega_0 t}$$

虚指数信号是周期信号：

$$x(t+kT) = e^{j\omega_0(t+kT)} = e^{j\omega_0 t} e^{j\omega_0 kT} = e^{j\omega_0 t} e^{j\frac{2\pi}{T}kT} = e^{j\omega_0 t} e^{j2k\pi} = e^{j\omega_0 t} = x(t)$$

其中周期为：

$$T = \frac{2\pi}{\omega_0}$$

Euler公式：

$$\cos(\omega t) = \frac{1}{2}(e^{j\omega t} + e^{-j\omega t}) \quad \sin(\omega t) = \frac{1}{2j}(e^{j\omega t} - e^{-j\omega t})$$

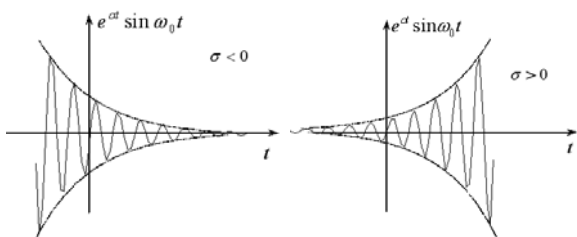
59

44

### 4. 复指数信号

$$x(t) = Ae^{st}, \quad s = \sigma + j\omega_0$$

$$x(t) = Ae^{\sigma t} e^{j\omega_0 t} = Ae^{\sigma t} \cos \omega_0 t + jAe^{\sigma t} \sin \omega_0 t$$



59

45

### 5. 抽样信号

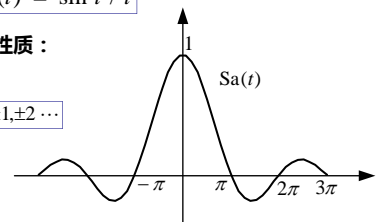
$$Sa(t) = \sin t / t$$

抽样函数具有以下性质：

$$Sa(0) = 1$$

$$Sa(k\pi) = 0, \quad k = \pm 1, \pm 2, \dots$$

$$\int_{-\infty}^{\infty} Sa(t) dt = \pi$$



与Sa(t)函数类似的是sinc(t)函数，其定义为：

$$\text{sinc}(t) = \sin(\pi t) / (\pi t)$$

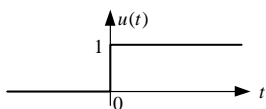
59

46

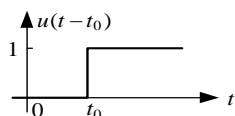
### 6. 阶跃信号

• 定义：

$$x(t) = \begin{cases} 1 & t > 0 \\ 0 & t < 0 \end{cases} = u(t)$$



$$u(t-t_0) = \begin{cases} 1 & t > t_0 \\ 0 & t < t_0 \end{cases}$$

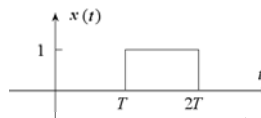


59

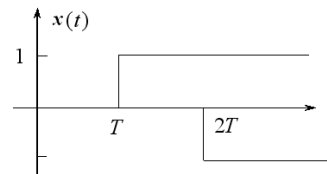
47

• 作用：

可表示任意的矩形脉冲信号



$$x(t) = u(t-T) - u(t-2T)$$



59

48



**可表示信号的时间范围**

59 49

**7. 斜坡信号**

• **定义：**

$$r(t) = \begin{cases} t & t \geq 0 \\ 0 & t < 0 \end{cases}$$

或  $r(t) = t \cdot u(t)$

• **与阶跃信号之间的关系：**

$$r(t) = \int_{-\infty}^t u(\tau) \cdot d\tau \quad \frac{dr(t)}{dt} = u(t)$$

59 50

**8. 冲激信号**

1) **冲激信号的引出**  
阶跃信号加在电容两端，流过电容的电流  $i(t) = Cdu(t)/dt$  可用冲激信号表示

2) **冲激信号的定义**  
**狄拉克定义式：**

$$\delta(t) = 0, t \neq 0$$

$$\int_{-\infty}^{+\infty} \delta(t) dt = 1$$

59 51

3) **冲激信号的图形表示**

• **说明：**  
• 冲激信号可延时至任意时刻  $t_0$ ，以符号  $\delta(t-t_0)$  表示

**定义式为：**  $\delta(t-t_0) = 0 \quad t \neq t_0$   
 $\int_{-\infty}^{\infty} \delta(t-t_0) dt = \int_{t_0-\Delta}^{t_0+\Delta} \delta(t-t_0) dt = 1$

**波形表示**

59 52

- **冲激信号具有强度，强度就是冲激信号对时间的定积分值**  
图中用括号注明，以区分信号的幅值
- **物理意义：作用时间极短，作用值很大的物理现象的数学模型**
- **冲激信号的作用：**  
表示其他任意信号      表示信号间断点的导数

---

4) **冲激信号的极限模型**

$$\delta(t) = \lim_{\Delta \rightarrow 0} x(t)$$

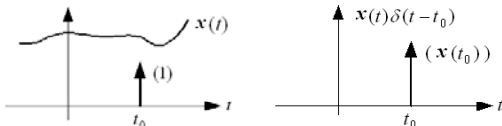
59 53

$$\delta(t) = \lim_{\Delta \rightarrow 0} f_{\Delta}(t) = \lim_{\Delta \rightarrow 0} g_{\Delta}(t) = \lim_{\Delta \rightarrow 0} h_{\Delta}(t)$$

59 54

### 5) 冲激信号的性质

- 筛选特性  $x(t)\delta(t-t_0) = x(t_0)\delta(t-t_0)$



- 取样特性  $\int_{-\infty}^{\infty} x(t)\delta(t-t_0)dt = x(t_0)$

证:

$$\int_{-\infty}^{\infty} x(t)\delta(t-t_0)dt = \int_{-\infty}^{\infty} x(t_0)\delta(t-t_0)dt = x(t_0) \int_{-\infty}^{\infty} \delta(t-t_0)dt = x(t_0)$$

59

55

- 展缩特性  $\delta(at) = \frac{1}{|a|}\delta(t)$

证: 任取一信号  $g(t)$ , 对于  $\int g(t)\delta(at)dt$

令  $x=at$ , 若  $a>0$ , 则有:

$$\int_{-\infty}^{\infty} g(t)\delta(at)dt = \int_{-\infty}^{\infty} g\left(\frac{x}{a}\right)\delta(x)\frac{dx}{a} = \frac{1}{a} \int_{-\infty}^{\infty} g\left(\frac{x}{a}\right)\delta(x)dx$$

若  $a<0$ , 则有:

$$\int_{-\infty}^{\infty} g(t)\delta(at)dt = \int_{\infty}^{-\infty} g\left(\frac{x}{a}\right)\delta(x)\frac{dx}{a} = \int_{\infty}^{-\infty} g\left(\frac{x}{a}\right)\delta(x)\frac{dx}{-a} = \frac{1}{-a} \int_{\infty}^{-\infty} g\left(\frac{x}{a}\right)\delta(x)dx$$

$$\therefore \int_{-\infty}^{\infty} g(x)\delta(ax)dx = \frac{1}{|a|} \int_{-\infty}^{\infty} g\left(\frac{x}{a}\right)\delta(x)dx = \frac{g(0)}{|a|}$$

$$\int_{-\infty}^{\infty} g(t)\frac{\delta(t)}{|a|}dt = \frac{1}{|a|} \int_{-\infty}^{\infty} g(t)\delta(t)dt = \frac{g(0)}{|a|} \Rightarrow \delta(at) = \frac{1}{|a|}\delta(t)$$

59

56

$$\delta(at) = \frac{1}{|a|}\delta(t)$$

$$a=-1 \Rightarrow \delta(t) = \delta(-t)$$

推论: 冲激信号是偶函数

- 与阶跃信号的关系

$$\int_{-\infty}^t \delta(\tau)d\tau = \begin{cases} 1 & t > 0 \\ 0 & t < 0 \end{cases} = u(t) \quad \frac{du(t)}{dt} = \delta(t)$$

### 6) 其他定义形式:

$$\delta(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{\pm j\omega t} d\omega, \quad \delta(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \cos \omega t d\omega$$

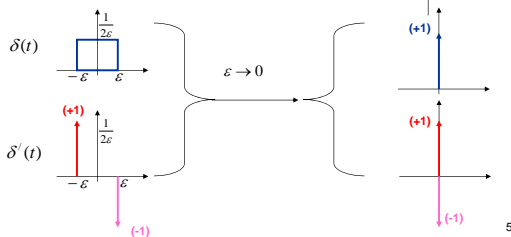
59

57

### 9. 冲激偶信号

- 定义:  $\frac{d\delta(t)}{dt} = \delta'(t)$  冲激偶信号图形表示

$$\lim_{\epsilon \rightarrow 0} x(t) = \frac{d\delta(t)}{dt} = \delta'(t)$$



59

58

- 性质:

$$\int_{-\infty}^{\infty} \delta'(t)dt = 0$$

$$\int_{-\infty}^t \delta'(\tau)d\tau = \delta(t)$$

$$\delta(t) = -t\delta'(t) \Rightarrow \delta'(t) = -\delta'(-t)$$

$$x(t)\delta'(t) = x(0)\delta'(t) - x'(0)\delta(t)$$

$$\int_{-\infty}^{\infty} x(t)\delta'(t)dt = -x'(0)$$

59

59